

Applied/Numerical Qualifier Solution: August 2009

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Problem 1. Consider the following finite element triple:

- $K =$ a rectangle with vertices $\{a^i\}$, $i = 1, 2, 3, 4$.
- $P = Q^3 = \text{span}\{x_1^i x_2^j; i, j = 0, \dots, 3\}$.
- $N = \{p(a^i), p_1(a^i), p_2(a^i), p_{12}(a^i), i = 1, 2, 3, 4\}$. (Here p_i denotes differentiation with respect to x_i).

a. Show that the above finite element is unisolvent.

Solution: For simplicity, we will work on the unit square \hat{K} with vertices

$$\{a^i\} = \{(0, 0), (1, 0), (1, 1), (0, 1)\}.$$

Working on the reference element is satisfactory due to the affine equivalence of finite elements. Recall the definition:

Definition: Let \mathcal{T}_h be a triangulation of $\Omega \subset \mathbb{R}^d$, $K \in \mathcal{T}_h$, P_K a finite dimensional vector space of functions defined on K , and Σ_K the space of linear forms mapping P_K to \mathbb{R} . Let further $(\hat{K}, \hat{P}_{\hat{K}}, \hat{\Sigma}_{\hat{K}})$, be a reference element. Then, the finite elements (K, P_K, Σ_K) , $K \in \mathcal{T}_h$, are said to be affine equivalent to the reference element $(\hat{K}, \hat{P}_{\hat{K}}, \hat{\Sigma}_{\hat{K}})$, if there exists an invertible affine mapping $F_K : \mathbb{R}^d \rightarrow \mathbb{R}^d$ such that for all $K \in \mathcal{T}_h$

$$K = F_K(\hat{K}) \tag{1}$$

$$P_K = \{p : K \rightarrow \mathbb{R} \mid p = \hat{p} \circ F_K^{-1}, \hat{p} \in \hat{P}_{\hat{K}}\} \tag{2}$$

$$\Sigma_K = \{\ell_i : P_K \rightarrow \mathbb{R} \mid \ell_i = \hat{\ell}_i \circ F_K^{-1}, \hat{\ell}_i \in \hat{\Sigma}_{\hat{K}}, 1 \leq i \leq n_K\}. \tag{3}$$

The idea is that the unisolvence on the reference element can be “transferred” to the physical elements. Let $f \in P$, then notice that for f restricted to the edge $[0, 1] \times \{x_2 = 0\}$, f is a third

degree polynomial. It can be seen that since f is zero at two points and f' is zero at two points, then f will be identically zero on that edge. Similarly, we can show that $f \equiv 0$ on the other three edges of our square. This implies that our function has the following form,

$$f(x_1, x_2) = x_1(x_1 - 1)x_2(x_2 - 1)\gamma(x_1, x_2),$$

where γ is a of the form $\gamma(x_1, x_2) = a + bx_1 + cx_2 + dx_1x_2$. Now consider the mixed partial derivative of f . Using product rule, we have,

$$\begin{aligned} f_{x_1x_2} &= (2x_1-1)(2x_2-1)\gamma + (2x_1-1)(x_2^2-x_2)\gamma_{x_2} \\ &\quad + (x_1^2-x_1)(2x_2-1)\gamma_{x_1} + (x_1^2-x_1)(x_2^2-x_2)\gamma_{x_1x_2}. \end{aligned}$$

Note that $(2x_1-1)(2x_2-1) \neq 0$ for our vertices $\{a^i\}$, but the other three terms will vanish for each a^i . Therefore we must have that $\gamma(a^i) = 0$ for $i = 1, 2, 3, 4$. But by definition of γ , if $\gamma = 0$ for four points then $\gamma \equiv 0$. Thus the finite element is unisolvent. ■

b. What do you need to do to check if the above element with a rectangular mesh results in a C^1 finite element space?

Solution: Let $p, q \in P$, with p defined on K_1 and q defined on K_2 where K_1 and K_2 are adjacent elements such that, p and q share the same values of the degrees of freedom on the interface between K_1 and K_2 . Then we must show that the “stiching together” of the functions p and q on the whole domain $K_1 \cup K_2$ results in a C^1 function. Specifically, we require that the function,

$$\phi(x, y) := \begin{cases} p(x, y) & \text{if } (x, y) \in K_1, \\ q(x, y) & \text{if } (x, y) \in K_2, \end{cases} \quad (4)$$

be a C^1 function. ■

c. Does the above element (with a rectangular mesh) result in a C^1 finite element space? (Explain your answer).

Proof: We follow the plan proposed in part b. Let $a^1 = (x_1, y_1)$ and $a^2 = (x_1, y_2)$ be the

points that form the line segment ℓ ; in particular, ℓ is a vertical line segment. Then we have,

$$\begin{aligned} p|_{\ell}(a^1) &= (q|_{\ell})(a^1), & p|_{\ell}(a^2) &= (q|_{\ell})(a^2) \\ (p_1)|_{\ell}(a^1) &= (q_1)|_{\ell}(a^1), & (p_1)|_{\ell}(a^2) &= (q_1)|_{\ell}(a^2) \\ (p_2)|_{\ell}(a^1) &= (q_2)|_{\ell}(a^1), & (p_2)|_{\ell}(a^2) &= (q_2)|_{\ell}(a^2) \\ (p_{12})|_{\ell}(a^1) &= (q_{12})|_{\ell}(a^1), & (p_{12})|_{\ell}(a^2) &= (q_{12})|_{\ell}(a^2) \end{aligned}$$

If we define,

$$\begin{aligned} f &= p|_{\ell} - q|_{\ell}, \\ g &= p_1|_{\ell} - q_1|_{\ell}. \end{aligned}$$

Then $f = f(y)$ is a cubic polynomial. Note that $f(y_2) = f(y_1) = f'(y_2) = f'(y_1) = 0$ which implies that f has 4 roots, hence $f \equiv 0$. By a similar reasoning, we can conclude that $g \equiv 0$. This implies that our function ϕ is continuous on $K_1 \cup K_2$. In addition, the partial derivatives exist and are continuous on all of $K_1 \cup K_2$. These results come directly from the shared degrees of freedom. Thus $\phi \in C^1(K_1 \cup K_2)$. A similar argument can be used for a horizontal line segment. This shows that the above element will result in a C^1 finite element space. ■

Problem 2. Consider the Neumann Problem:

$$-\Delta u = f \quad \text{in } \Omega \tag{5}$$

$$\frac{\partial u}{\partial n} = g \quad \text{on } \partial\Omega. \tag{6}$$

Here Ω is a bounded domain in \mathbb{R}^2 and f and g are suitably smooth.

a. Derive a weak form of the above problem using a test function in $H^1(\Omega)$.

Proof: Multiplying $-\Delta u = f$ by a test function $v \in H^1(\Omega)$ and integrating over Ω , we have,

$$\begin{aligned} - \int_{\Omega} \Delta u v \, d\mathbf{x} &= \int_{\Omega} \nabla u \cdot \nabla v \, d\mathbf{x} - \int_{\partial\Omega} \frac{\partial u}{\partial n} v \, ds \\ &= \int_{\Omega} \nabla u \cdot \nabla v \, d\mathbf{x} - \int_{\partial\Omega} g v \, ds. \end{aligned}$$

Adding the boundary term to the other side, we have,

$$a(u, v) := \int_{\Omega} \nabla u \cdot \nabla v \, d\mathbf{x} = \int_{\Omega} f v \, d\mathbf{x} + \int_{\partial\Omega} g v \, ds =: L(v).$$

So our weak formulation becomes: find $u \in H^1(\Omega)$ such that $a(u, v) = L(v)$ for all $v \in H^1(\Omega)$.

b. Discuss when the weak form of Part a. has a solution and if it is unique.

Proof: First note that a and F can both be shown to be continuous (for F you need to invoke the trace lemma). Then from the Lax-Milgram lemma a solution exists and is unique if a is $H^1(\Omega)$ -elliptic (coercive). However, in our case $H^1(\Omega)$ -ellipticity will not hold since we do not have a Poincaré inequality (or any other means to bound $a(u, u)$ below). Note that $v = \text{constant}$ is a permissible function in our test function space. So set $v = 1$ and we have

$$\int_{\Omega} f \, d\mathbf{x} + \int_{\partial\Omega} g \, ds = 0.$$

This is called the *compatibility condition* or *solvability condition* and is necessary for existence of the solution.

Next, note that if u is a solution then $u + c$ is also a solution for $c \in \mathbb{R}$. In order to guarantee a *unique* solution we can impose the condition,

$$\int_{\Omega} u \, d\mathbf{x} = 0.$$

This assumption allows us to derive a Poincaré inequality, which allows us apply the Lax-Milgram lemma. ■

c. Describe a variational formulation of (5) in terms of an appropriate Hilbert space V . Be sure to explicitly define V .

Proof: The variational formulation will be: find $u \in V := \{v \in H^1(\Omega) : \int_{\Omega} v \, d\mathbf{x} = 0\}$ such that $a(u, v) = L(v)$ for all $v \in V$. Note we cannot impose both $\partial u / \partial n = 0$ and $u = 0$ (Cauchy boundary condition) on $\partial\Omega$ since then the problem becomes ill-posed.

d. Prove coercivity of the form of Part a. on the V of Part c. when $\Omega = (0, 1)^2$.

Proof: In order to prove coercivity, we use a Poincaré type inequality. Specifically, we will prove

$$\|u\|_{L^2(\Omega)}^2 \leq C \left[\left(\int_{\Omega} u \, dx \right)^2 + \|\nabla u\|_{L^2(\Omega)}^2 \right] = C \|\nabla u\|_{L^2(\Omega)}^2,$$

for some constant $C > 0$. We will start by proving the result for a smooth function $\phi \in C^\infty(\Omega) \cap V$. Note we have the following identity,

$$\phi(x_2, y_2) - \phi(x_1, y_1) = \int_{x_1}^{x_2} \frac{\partial \phi}{\partial x}(x, y_2) \, dx + \int_{y_1}^{y_2} \frac{\partial \phi}{\partial y}(x_1, y) \, dy.$$

Squaring both sides, we have,

$$\begin{aligned} \phi^2(x_2, y_2) - 2\phi(x_2, y_2)\phi(x_1, y_1) + \phi^2(x_1, y_1) &= \left(\int_{x_1}^{x_2} \frac{\partial \phi}{\partial x}(x, y_2) \, dx + \int_{y_1}^{y_2} \frac{\partial \phi}{\partial y}(x_1, y) \, dy \right)^2 \\ &\leq 2 \left(\int_{x_1}^{x_2} \frac{\partial \phi}{\partial x}(x, y_2) \, dx \right)^2 + 2 \left(\int_{y_1}^{y_2} \frac{\partial \phi}{\partial y}(x_1, y) \, dy \right)^2 \\ &\leq 2 \|\nabla \phi\|_{L^2(\Omega)}^2. \end{aligned}$$

Now we integrate both sides of the inequality from 0 to 1 for x_1 , x_2 , y_1 , and y_2 (four different variables). Doing so will give us,

$$2\|\phi\|_{L^2(\Omega)}^2 - 2 \left(\int_{\Omega} \phi \, dx \, dy \right)^2 \leq 2\|\nabla \phi\|_{L^2(\Omega)}^2.$$

Thus the result follows where $C = 1$. Now to prove for $u \in H^1(\Omega)$, we use the fact that

$$H^1(\Omega) = \overline{C^\infty(\Omega)}^{\|\cdot\|_{H^1(\Omega)}}.$$

So for $\varepsilon > 0$, we can find $\phi \in C^\infty(\Omega) \cap V$ such that $\|u - \phi\|_{H^1(\Omega)}^2 < \varepsilon/16$. Now consider,

$$\begin{aligned} \|u\|_{L^2(\Omega)}^2 &\leq \|u\|_{H^1(\Omega)}^2 \\ &\leq (\|u - \phi\|_{H^1(\Omega)} + \|\phi\|_{H^1(\Omega)})^2 \\ &\leq 2\|u - \phi\|_{H^1(\Omega)}^2 + 2\|\phi\|_{H^1(\Omega)}^2 \\ &\leq \frac{\varepsilon}{8} + 2\|\phi\|_{L^2(\Omega)}^2 + 2\|\phi\|_{H^1(\Omega)}^2 \\ &\leq \frac{\varepsilon}{8} + 4\|\phi\|_{H^1(\Omega)}^2 \\ &\leq \frac{\varepsilon}{8} + 8\|u - \phi\|_{H^1(\Omega)}^2 + 8\|u\|_{H^1(\Omega)}^2. \end{aligned}$$

Then note that $\|u - \phi\|_{H^1(\Omega)}^2 \leq \|u - \phi\|_{H^1(\Omega)}^2 < \varepsilon/16$. Thus we have,

$$\|u\|_{L^2(\Omega)}^2 \leq \frac{\varepsilon}{8} + \frac{\varepsilon}{2} + 8\|u\|_{H^1(\Omega)}^2 < \varepsilon + 8\|u\|_{H^1(\Omega)}^2.$$

Since this holds for any $\varepsilon > 0$, we thus have $\|u\|_{L^2(\Omega)} \leq C\|\nabla u\|_{H^1(\Omega)}$. Since we have a Poincaré inequality, the coercivity proof follows in the typical way. ■

Problem 3. Let $\Omega_e = \{x \in \mathbb{R}^2 : \|x\| > 1\}$. Show that the Poincarè inequality does not hold in $H_0^1(\Omega_e)$, i.e., there does not exist a constant $c > 0$ satisfying

$$c\|u\|_{L^2(\Omega_e)}^2 \leq \int_{\Omega_e} \|\nabla u\|^2 dx \quad \text{for all } u \in H_0^1(\Omega_e). \quad (7)$$

The space $H_0^1(\Omega_e)$ is the completion of $C_0^\infty(\Omega_e)$ in the norm

$$\|v\|_{H^1(\Omega_e)} = \left(\|v\|_{L^2(\Omega_e)}^2 + \|\nabla v\|_{(L^2(\Omega_e))^2}^2 \right)^{1/2}. \quad (8)$$

(Hint: Consider dilating a fixed function.)

Proof: Consider the rotationally symmetric bump function,

$$\phi(r, \theta) := \phi(r) := \begin{cases} \exp(\frac{-1}{r(1-r)}) & \text{for } 0 < r < 1 \\ 0 & \text{for } r \geq 1, \end{cases}$$

which is only a function of r . Consider the sequence of bump functions defined by,

$$\phi_n(r) := \phi\left(\frac{r-1}{n}\right).$$

Note that $\text{supp}(\phi_n) = [1, n+1]$. Then consider the L^2 norm of ϕ_n ,

$$\begin{aligned} \|\phi_n\|_{L^2(\Omega_e)}^2 &= \int_0^{2\pi} \int_1^\infty \phi_n^2(r) r dr d\theta \\ &= 2\pi \int_1^{n+1} \phi^2\left(\frac{r-1}{n}\right) r dr \\ &= 2\pi \int_0^1 \phi^2(s) n(ns+1) ds \\ &\geq 2\pi n^2 \int_0^1 s \phi^2(s) ds \\ &= C_0 n^2. \end{aligned}$$

Thus $\lim_{n \rightarrow \infty} \|\phi_n\|_{L^2(\Omega_e)} = \infty$. But consider the H^1 -semi norm of ϕ_n .

$$\begin{aligned} |\phi_n|_{H^1(\Omega_e)}^2 &= \int_0^{2\pi} \int_1^\infty |\phi_n'(r)|^2 r dr d\theta \\ &= 2\pi \int_1^{n+1} \left| \frac{\partial}{\partial r} \phi\left(\frac{r-1}{n}\right) \right|^2 r dr \\ &= 2\pi \int_0^1 \left| \frac{1}{n} \frac{\partial}{\partial s} \phi(s) \right|^2 n(ns+1) ds \\ &\leq 2\pi \frac{n(n+1)}{n^2} \int_0^1 |\phi'(s)|^2 ds \\ &\leq C_1 \end{aligned}$$

Thus since the H^1 -semi norm is bounded for all n and the L^2 norm blows up to infinity, this implies that we cannot have a Poincarè inequality on this domain. ■